Numerical Integration

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Extra

Quadrature techniques

$$I = \int_a^b f(x) dx \approx \sum_{i=1}^n w_i f(x_i) = \sum_{i=1}^n w_i f_i$$

- Nodes: *x_i*
- Weights: w_i

Extra

Quadrature techniques

$$I = \int_{a}^{b} f(x) dx \approx \sum_{i=1}^{n} w_{i} f(x_{i})$$

Two versions:

- Newton Cotes:
 - equidistant nodes & "best" choice for the weights w_i
- Gaussian Quadrature:
 - "best" choice for both nodes and weights

Monte Carlo techniques

- pseudo:
 - implemetable version of true Monte Carlo
- quasi:
 - looks like Monte Carlo, but is something different
 - name should have been chosen better

Power

- Newton-Cotes: With n nodes you get
 - exact answer if f is $(n-1)^{\rm th}\text{-order polynomial}$
 - accurate answer f is close to a n^{th} -order polynomial
- Gaussian: With n nodes you get
 - exact answer if f is $(2n-1)^{\text{th}}$ -order polynomial
 - accurate answer f is close to a $\left(2n-1\right)^{\mathrm{th}}$ -order polynomial

Power

- (Pseudo) Monte Carlo: accuracy requires lots of draws
- Quasi Monte Carlo: definitely better than (pseudo) Monte Carlo and will dominate quadrature methods for higher-dimensional problems

Idea behind Newton-Cotes

 function values at n nodes ⇒ you can fit a (n − 1)th-order polynomial & integrate the approximating polynomial

$$\int_{a}^{b} f(x) dx \approx \int_{a}^{b} P_{2}(x) dx$$

- It turns out that this can be standardized
 - see Section "extra" for derivation

Simpson for one interval (3 nodes)

$$\int_{a}^{b} f(x) dx \approx \left(\frac{1}{3}f_{0} + \frac{4}{3}f_{1} + \frac{1}{3}f_{2}\right) h$$

Simpson for multiple intervals (n+1 nodes)

$$\int_{a}^{b} f(x)dx \approx \left(\frac{1}{3}f_{0} + \frac{4}{3}f_{1} + \frac{1}{3}f_{2}\right)h \\ + \left(\frac{1}{3}f_{2} + \frac{4}{3}f_{3} + \frac{1}{3}f_{4}\right)h \\ + \cdots \\ + \left(\frac{1}{3}f_{n-2} + \frac{4}{3}f_{n-1} + \frac{1}{3}f_{n}\right)h \\ = \left(\frac{1}{3}f_{0} + \frac{4}{3}f_{1} + \frac{2}{3}f_{2} + \frac{4}{3}f_{3} + \frac{2}{3}f_{4} + \cdots + \frac{2}{3}f_{n-2} + \frac{4}{3}f_{n-1} + \frac{1}{3}f_{n}\right)h$$

Gaussian quadrature

- Could we do better? That is, get better accuracy with same amount of nodes?
- Answer: Yes, if you are smart about choosing the nodes
 - This is Gaussian quadrature

Gauss-Legendre quadrature

- Let [a,b] be [-1,1]
 - can always be accomplished by scaling
- Quadrature

$$\int_{-1}^{1} f(x) dx \approx \sum_{i=1}^{n} \omega_i f(\zeta_i).$$

- **Goal:** Get exact answer if f(x) is a polynomial of order 2n 1
- That is with 5 nodes you get exact answer even if f(x) is a $9^{\rm th}$ -order polynomial

Implementing Gauss-Legendre quadrature

• Get n nodes and n weights from a computer program

•
$$\zeta_i$$
, $i = 1, \cdots, n, \omega_i, i = 1, \cdots, n$

- Calculate the function values at the *n* nodes, $f_i \ i = 1, \cdots, n$
- Answer is equal to

$$\sum_{i=1}^n \omega_i f_i$$

- Anybody could do this
- How does the computer get the nodes and weights?

2n equations for nodes and weights

• To get right answer for f(x) = 1

$$\int_{-1}^{1} 1 dx = \sum_{i=1}^{n} \omega_i 1$$

• To get right answer for f(x) = x

$$\int_{-1}^{1} x dx = \sum_{i=1}^{n} \omega_i \zeta_i$$

• To get right answer for $f(x) = x^2$

$$\int_{-1}^{1} x^2 dx = \sum_{i=1}^{n} \omega_i \zeta_i^2$$

etc

2n equations for nodes and weights

• To get right answer for $f(x) = x^j$ for $j = 0, \cdots, 2n-1$

$$\int_{-1}^{1} x^{j} dx = \sum_{i=1}^{n} \omega_{i} \zeta_{i}^{j} \quad j = 0, 1, \cdots, 2n-1$$

• This is a system of 2n equations in 2n unknowns.

What has been accomplished so far?

• By construction we get right answer for

$$f(x) = 1, f(x) = x, \dots, f(x) = x^{2n-1}$$

• But this is enough to get right answer for any polynomial of order 2n - 1

$$f(x) = \sum_{i=0}^{2n-1} a_i x^i$$

• Why?

Gauss-Hermite Quadrature

• Suppose we want to approximate

$$\int_{-\infty}^{\infty} f(x)e^{-x^2}dx \text{ with } \sum_{i=1}^{n} \omega_i f(\zeta_i)$$

• The function e^{-x^2} is the *weighting function*, it is not used in the approximation but is captured by the ω_i coefficients

Gauss-Hermite Quadrature

• We can use the same procedure to find the weights and the nodes, that is we solve them from the system:

$$\int_{-\infty}^{\infty} x^{j} e^{-x^{2}} dx = \sum_{i=1}^{n} \omega_{i} \zeta_{i}^{j} \text{ for } j = 0, 1, \cdots, 2n-1$$

• Note that $e^{-\zeta_i^2}$ is *not* on the right-hand side

Implementing Gauss-Hermite Quadrature

- Get n nodes, ζ_i, i = 1, · · · , n, and n weights, ω_i, i = 1, · · · , n, from a computer program
- Calculate the function values at the n nodes, $f_i \ i = 1, \cdots, n$
- Answer is equal to

$$\sum_{i=1}^n \omega_i f_i$$

Extra

Expectation of Normally distributed variable

• How to calculate

$$\mathsf{E}\left[h(y)
ight]$$
 with $y \sim N(\mu, \sigma^2)$

• That is, we have to calculate

$$\int_{-\infty}^{\infty} \frac{1}{\sigma\sqrt{2\pi}} h(y) \exp\left(-\frac{(y-\mu)^2}{2\sigma^2}\right) dy$$

• Unfortunately, this does not exactly fit the Hermite weighting function, but a change in variable will do the trick

Gaussian quadrature

Extra

Change of variables

• If $y = \phi(x)$ then

$$\int_a^b g(y)dy = \int_{\phi^{-1}(a)}^{\phi^{-1}(b)} g(\phi(x))\phi'(x)dx$$

• Note the Jacobian is added

Gaussian quadrature

Extra

Change of variables

The transformation we use here is

$$x = rac{y-\mu}{\sigma\sqrt{2}}$$
 or $y = \sigma\sqrt{2}x + \mu$

Change of variables

$$\mathsf{E}\left[h(y)\right] = \int_{-\infty}^{\infty} \frac{1}{\sigma\sqrt{2\pi}} h(y) \exp\left(-\frac{(y-\mu)^2}{2\sigma^2}\right) dy$$

$$= \int_{-\infty}^{\infty} \frac{1}{\sigma\sqrt{2\pi}} h(\sqrt{2}\sigma x + \mu) \exp\left(-x^2\right) \sigma\sqrt{2} dx$$

$$=\int_{-\infty}^{\infty}\frac{1}{\sqrt{\pi}}h(\sqrt{2}\sigma x+\mu)\exp\left(-x^{2}\right)dx$$

What to do in practice?

- Obtains *n* Gauss-Hermite quadrature weights and nodes using a numerical algorithm.
- Calculate the approximation using

$$\mathsf{E}\left[h(y)\right] \approx \sum_{i=1}^{n} \frac{1}{\sqrt{\pi}} \omega_i^{GH} h\left(\sqrt{2}\sigma \zeta_i^{GH} + \mu\right)$$

- Do not forget to divide by $\sqrt{\pi}!$
- Is this amazingly simple or what?

Gaussian quadrature

Extra

Extra material

- Derivation Simpson formula
- Monte Carlo integration

Lagrange interpolation

Let

$$L_{i}(x) = \frac{(x - x_{0}) \cdots (x - x_{i-1})(x - x_{i+1}) \cdots (x - x_{n})}{(x_{i} - x_{0}) \cdots (x_{i} - x_{i-1})(x_{i} - x_{i+1}) \cdots (x_{i} - x_{n})}$$

$$f(x) \approx f_0 L_0(x) + \cdots + f_n L_n(x).$$

- What is the right-hand side?
- Do I have a perfect fit at the n + 1 nodes?

Simpson: 2nd-order Newton-Cotes

•
$$x_0 = a$$
, $x_1 = (a+b)/2$, $x_2 = b$, or

•
$$x_1 = x_0 + h$$
, $x_2 = x_0 + 2h$

Using the Lagrange way of writing the 2nd-order polynomial, we get

$$\int_{a}^{b} f(x)dx \approx \int_{a}^{b} f_{0}L_{0}(x) + f_{1}L_{1}(x) + f_{2}L_{2}(x)$$

= $f_{0}\int_{a}^{b} L_{0}(x)dx + f_{1}\int_{a}^{b} L_{1}(x)dx + f_{2}\int_{a}^{b} L_{2}(x)dx$

Amazing algebra

$$\int_a^b L_0(x)dx = \frac{1}{3}h$$
$$\int_a^b L_1(x)dx = \frac{4}{3}h$$
$$\int_a^b L_2(x)dx = \frac{1}{3}h$$

- Why amazing?
 - formula only depends on h, not on values x_i and f_i
- Combining gives

$$\int_{a}^{b} f(x) dx \approx \int_{a}^{b} P_{2}(x) dx = \left(\frac{1}{3}f_{0} + \frac{4}{3}f_{1} + \frac{1}{3}f_{2}\right) h.$$

True and pseudo Monte Carlo

To calculate an expectation

- Let x be a random variable with CDF F(x)
- Monte Carlo integration:

$$\int_{a}^{b} h(x) dF(x) \approx \frac{\sum_{t=1}^{T} h(x_t)}{T},$$

• Use random number generator to implement this in practice

True and pseudo Monte Carlo

What if integral is not an expectation

$$\int_a^b h(x)dx = (b-a)\int_a^b h(x)f_{\mathsf{ab}}(x)dx,$$

where f_{ab} is the density of a random variable with a uniform distribution over [a, b], that is, $f_{ab} = (b - a)^{-1}$. Thus, one could approximate the integral with

$$\int_{a}^{b} h(x) dx \approx (b-a) \frac{\sum_{t=1}^{T} h(x_t)}{T},$$

where x_t is generated using a random number generator for a variable that is uniform on [a, b].

Quasi Monte Carlo

- Monte Carlo integration has very slow convergence properties
- In higher dimensional problems, however, it does better than quadrature (it seems to avoid the curse of dimensionality)
- But why? Pseudo MC is simply a deterministic way to go through the state space
- Quasi MC takes that idea and improves upon it

Quasi Monte Carlo

- Idea: Fill the space in an efficient way
- Equidistributed series: A scalar sequence $\{x_t\}_{t=1}^T$ is equidistributed over [a, b] iff

$$\lim_{T \longrightarrow \infty} \frac{b-a}{T} \sum_{t=1}^{T} f(x_t) = \int_a^b f(x) dx$$

for all Rieman-integrable f(x).

• Equidistributed takes the place of uniform

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Quasi Monte Carlo

- Examples
 - $\xi, 2\xi, 3\xi, 4\xi, \cdots$ is equidistributed modulo 1 for any irrational number ξ .¹
 - The sequence of prime numbers multiplied by an irrational number (2ξ, 3ξ, 5ξ, 7ξ, · · ·)

 $^{{}^{1}}Frac(x)$ (or x Modulo 1) means that we subtract the largest integer that is less than x. For example, frac(3.564) = 0.564.

Extra

Multidimensional

For a *d*-dimensional problem, an equidistributed sequence $\{x_t\}_{t=1}^T \subset D \subset R^d$ satisfies

$$\lim_{T \longrightarrow \infty} \frac{\mu(D)}{T} \sum_{t=1}^{T} f(x_t) = \int_D f(x) dx,$$

where $\mu(D)$ is the Lebesque measure of D.

Multidimensional equidistributed vectors

Examples for the *d*-dimensional unit hypercube:

Weyl:

$$x_t = (t\sqrt{p_1}, t\sqrt{p_2}, \cdots, t\sqrt{p_d}) ext{ modulo 1,}$$

where p_i is the i^{th} positive prime number.

Neiderreiter:

$$x_t = (t2^{1/(d+1)}, t2^{2/(d+1)}, \cdots, t2^{d/(d+1)})$$
 modulo 1



References

- Den Haan, W.J., Numerical Integration
- Most text books on numerical methods will have a chapter on this topic